



RESEARCH ARTICLE

STUDY ON EINSTEIN- KAEHLERIAN DECOMPOSABLE RECURRENT SPACE OF FIRST ORDER

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ABSTRACT

Takano [1967] have studied decomposition of curvature tensor in a recurrent space. Sinha and Singh [1970] have been studied and defined decomposition of recurrent curvature tensor field in a Finsler space. Singh and Negi studied decomposition of recurrent curvature tensor field in a Kaehlerian space. Negi and Rawat [1995] have studied decomposition of recurrent curvature tensor field in Kaehlerian space. Rawat and Silswal [2007] studied and defined decomposition of recurrent curvature tensor fields in a Tachibana space. Further, Rawat and Kunwar Singh [2008] studied the decomposition of curvature tensor field in Kaehlerian recurrent space of first order. In the present paper, we have studied the decomposition of curvature tensor fields R_{ijk}^h in terms of two non-zero vectors and a tensor field in Einstein-Kaehlerian recurrent space of first order and several theorem have been established and proved.

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INTRODUCTION

An n (=2m) dimensional Kaehlerian space K_n is a Riemannian space, which admits a tensor field F_i^h satisfying the conditions

$$F_i^h F_h^j = -\delta_i^j \tag{1.1}$$

$$F_{ij} = -F_{ji}, (F_{ij} = F_i^a g_{aj}) \tag{1.2}$$

$$\text{and } F_{ij}^h = 0, \tag{1.3}$$

Where the (.) followed by an index denotes the operation of covariant differentiation with respect to the metric tensor g_{ij} of the Riemannian space.

The Riemannian curvature tensor R_{ijk}^h is given by

$$R_{ijk}^h = \partial_i \left\{ \begin{matrix} h \\ j k \end{matrix} \right\} - \partial_j \left\{ \begin{matrix} h \\ i k \end{matrix} \right\} + \left\{ \begin{matrix} h \\ i \alpha \end{matrix} \right\} \left\{ \begin{matrix} \alpha \\ j k \end{matrix} \right\} - \left\{ \begin{matrix} h \\ j \alpha \end{matrix} \right\} \left\{ \begin{matrix} \alpha \\ i k \end{matrix} \right\},$$

The Ricci tensor and the scalar Curvature tensor are respectively given by

$$R_{ij} = R_{aij}^a \text{ and } R = g_{ij} R_{ij} \tag{1.4}$$

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It is well known that these tensors satisfies the following identities

$$R_{ijk}^a = R_{jki} - R_{ikj} \quad (1.5)$$

$$R_{,i} = 2R_{i,a}^a \quad (1.6)$$

$$F_i^a R_{aj} = - R_{ia} F_j^a \quad (1.7)$$

and

$$F_i^a R_a^i = R_i^a F_a^i \quad (1.8)$$

The holomorphically projective curvature tensor P_{ijk}^h is defined by

$$P_{ijk}^h = R_{ijk}^h + \frac{1}{(n+2)} (R_{ik} \delta_j^h - R_{jk} \delta_i^h + S_{ik} F_j^h - S_{jk} F_i^h + 2S_{ij} F_k^h) \quad (1.9)$$

where $S_{ij} = F_i^a R_{aj}$

Let us suppose that a Kaehlerian space is Einstein one, and then the Ricci tensor satisfies

$$R_{ij} = \frac{R}{n} g_{ij}, R_{,a} = 0$$

from which, we obtain

$$R_{ij,a} = 0, S_{ij,a} = 0$$

$$\text{and } S_{ij} = \frac{R}{n} F_{ij}$$

The Bianchi identity for Einstein-Kaehlerian space are given by

$$R_{ijk}^h + R_{jki}^h + R_{kij}^h = 0 \quad (1.10)$$

and

$$R_{ijk,a}^h + R_{ika,j}^h + R_{iaj,k}^h = 0 \quad (1.11)$$

The Commutative formulae for the Curvature tensor fields are given as follows

$$T_{,jk}^i - T_{,kj}^i = T^a R_{ajk}^i \quad (1.12)$$

$$T_{i,ml}^h - T_{i,lm}^h = T_i^a R_{aml}^h - T_a^h R_{iml}^a \quad (1.13)$$

A Einstein-Kaehlerian space is said to be Einstein-Kaehlerian recurrent space of first order, if its curvature tensor field satisfy the condition

$$R_{ijk,a}^h = \lambda_a R_{ijk}^h \quad (1.14)$$

where λ_a is a non - zero vector and is known as recurrence vector field. The following relations follow immediately from equation (1.14),

$$R_{ij,a} = \lambda_a R_{ij} \quad (1.15)$$

and

$$R_{,a} = \lambda_a R \quad (1.16)$$

2. Decomposition of Curvature Tensor Field R_{ijk}^h

We Consider the decomposition of recurrent curvature tensor field R_{ijk}^h in the following form

$$R_{ijk}^h = V_i^h \Phi_j \Psi_k \quad (2.1)$$

where the non - zero tensor field V_i^h and vector Φ_j, Ψ_k are such that

$$\lambda_h V_i^h = P_i \quad (2.2)$$

Theorem (2.1) : Under the decomposition (2.1), the Bianchi identities for R_{ijk}^h takes the forms

$$P_i \Phi_j \Psi_k + P_j \Phi_k \Psi_i + P_k \Phi_i \Psi_j = 0 \quad (2.3)$$

and

$$\lambda_a \Phi_j \Psi_k + \lambda_j \Phi_k \Psi_a + \lambda_k \Phi_a \Psi_j = 0 \quad (2.4)$$

Proof:- From Equations (1.10) and (2.1), we have

$$V_i^h \Phi_j \Psi_k + V_j^h \Phi_k \Psi_i + V_k^h \Phi_i \Psi_j = 0 \quad (2.5)$$

Multiplying (2.5), by λ_h , and using (2.2), we get relation (2.3)

$$P_i \Phi_j \Psi_k + P_j \Phi_k \Psi_i + P_k \Phi_i \Psi_j = 0$$

From Equations (1.11), (1.14) and (2.1), we have

$$V_i^h [\lambda_a \Phi_j \Psi_k + \lambda_j \Phi_k \Psi_a + \lambda_k \Phi_a \Psi_j] = 0 \quad (2.6)$$

Multiplying (2.6) by λ_h and using (2.2), we get relation (2.4).

Theorem (2.2) : Under the decomposition (2.1), the tensor field R_{ijk}^h , R_{ij} and vectors Φ_j, Ψ_k satisfies the relations

$$\lambda_a R_{ijk}^a = \lambda_i R_{jk} - \lambda_j R_{ik} = P_i \Phi_j \Psi_k \quad (2.7)$$

Proof : With the help of Equations (1.5), (1.14) and (1.15), we have

$$\lambda_a R_{ijk}^a = \lambda_i R_{jk} - \lambda_j R_{ik} \quad (2.8)$$

Multiplying (2.1) by λ_h , and using relation (2.2), we get

$$\lambda_h R_{ijk}^h = P_i \Phi_j \Psi_k \quad (2.9)$$

in view of (2.8) and (2.9), we get the required relation (2.7).

Theorem (2.3) : Under the decomposition (2.1), the quantities λ_a and V_i^h behave like the recurrent vector and tensor field. The recurrent form of these quantities are given by

$$\lambda_{a,m} = \mu_m \lambda_a \quad (2.10)$$

$$V_{i,m}^h = \mu_m V_i^h \quad (2.11)$$

Proof : Differentiating (2.7), covariantly with respect to x^m , and using (2.1) and (2.7), we obtain

$$\lambda_{a,m} V_i^a \Phi_j \Psi_k = \lambda_{i,m} R_{jk} - \lambda_{j,m} R_{ik} \quad (2.12)$$

Multiplying (2.12) by λ_a and using (2.1) and (2.8), we get

$$\lambda_{a,m} (\lambda_i R_{jk} - \lambda_j R_{ik}) = \lambda_a (\lambda_{i,m} R_{jk} - \lambda_{j,m} R_{ik}) \quad (2.13)$$

Now, multiplying (2.13) by λ_h , we have

$$\lambda_{a,m} (\lambda_i R_{jk} - \lambda_j R_{ik}) \lambda_h = \lambda_a \lambda_h (\lambda_{i,m} R_{jk} - \lambda_{j,m} R_{ik}) \quad (2.14)$$

Since the expression of right hand side of the above equation is symmetric in a and h, therefore

$$\lambda_{a,m}\lambda_h = \lambda_{h,m}\lambda_a, \quad (2.15)$$

Provided that $\lambda_i R_{jk} - \lambda_j R_{ik} \neq 0$

The vector field λ_a being non-zero, we can have a proportional vector μ_m such that

$$\lambda_{a,m} = \mu_m \lambda_a \quad (2.16)$$

Further, differentiating (2.2) w.r to x^m and using (2.16), we get

$$\lambda_h V_{i,m}^h = P_{i,m} - \mu_m P_i \quad (2.17)$$

from the above equation, it is obvious that

$$\lambda_h V_{i,m}^h = \lambda_a V_{i,m}^a \quad (2.18)$$

Since λ_a is a non-zero recurrence vector field, we can get a proportional vector field μ_m such that

$$V_{i,m}^h = \mu_m V_i^h$$

which complete the proof.

Theorem (2.4) : Under the decomposition (2.1), the vector field P_i, Φ_j, Ψ_k behave like recurrent vectors and their recurrent form are given respectively by

$$P_{i,m} = 2 \mu_m P_i \quad (2.19)$$

and

$$(\lambda_m - \mu_m) \Phi_j \Psi_k = \Phi_{j,m} \Psi_k + \Phi_j \Psi_{k,m} \quad (2.20)$$

Proof Differentiating (2.2) covariantly w.r. to x^m , and using equation (2.2), (2.10) and (2.11), we obtain the required result (2.19). Further, differentiating equation (2.1) covariantly w.r. to x^m , and using equation (1.14), (2.1) and (2.11), we get the required recurrent form (2.20).

Theorem (2.5) : Under the decomposition (2.1), the curvature tensor and holomorphically projective curvature tensor are equal if

$$\Phi_k \Psi_i \{ (P_i \delta_j^h - P_j \delta_i^h) + P_a (F_i^a F_j^h - F_j^a F_i^h) \} + 2 P_a \Phi_j \Psi_i F_i^a F_k^h = 0 \quad (2.21)$$

Proof: The equation (1.9), may be written in the form

$$P_{ijk}^h = R_{ijk}^h + D_{ijk}^h \quad (2.22)$$

where

$$D_{ijk}^h = \frac{1}{n+2} (R_{ik} \delta_j^h - R_{jk} \delta_i^h + S_{ik} F_j^h - S_{jk} F_i^h + 2 S_{ij} F_k^h) \quad (2.23)$$

Contracting indices h and k in (2.1), we obtain

$$R_{ij} = V_i^l \Phi_j \Psi_l \quad (2.24)$$

In view of equation (2.24), we have

$$S_{ij} = F_i^a \Phi_j \Psi_l V_a^l \quad (2.25)$$

Making use of (2.24) and (2.25) in equation (2.22), we get

$$D_{ijk}^h = \frac{1}{n+2} [\Phi_k \Psi_l \{ (V_i^l \delta_j^h - V_j^l \delta_i^h) + V_a^l (F_i^a F_j^h - F_j^a F_i^h) \} + 2 \Phi_j \Psi_l F_i^a F_k^h V_a^l] \quad (2.26)$$

In view of (2.23), it is clear that

$$P_{ijk}^h = R_{ijk}^h \text{ iff } D_{ijk}^h = 0, \text{ which in view of equation (2.26) gives}$$

$$[\Phi_k \Psi_l \{(V_i^l \delta_j^h - V_j^l \delta_i^h) + V_a^l (F_i^a F_j^h - F_j^a F_i^h)\} + 2 \Phi_j \Psi_l F_i^a F_k^h V_a^l] = 0 \quad (2.27)$$

Multiplying (2.27) by λ_l and using (2.2), we obtain the required condition (2.21).

Theorem (2.6) : Under the decomposition (2.1), the scalar curvature R , satisfy the relation.

$$\lambda_k R = g^{ij} P_i \Phi_j \Psi_l \quad (2.28)$$

Proof : Contracting indices h and k in (2.1), we get

$$R_{ij} = V_i^l \Phi_j \Psi_l \quad (2.29)$$

Multiplying (2.29) by g^{ij} both sides, we get

$$g^{ij} R_{ij} = g^{ij} V_i^l \Phi_j \Psi_l \quad (2.30)$$

in view of Equation (1.4), the above equation reduces to

$$R = g^{ij} V_i^l \Phi_j \Psi_l \quad (2.31)$$

Now multiplying (2.31) by λ_l and using (2.2), we get

$$\lambda_l R = g^{ij} P_i \Phi_j \Psi_l$$

which complete the proof of the theorem.

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